

The Expansion of the Inverse of the Mutual Distance Between Two Planets Raised to Any Real Integer Using Taylor Theorem and Elliptic Expansions Part I: Theoretical Results

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We introduce in this part the method to obtain the literal expansion of the mutual distance between two planets of the solar system raised to any negative real integer.

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1. Introduction

The subject of the expansion of planetary disturbing function is extensive. It plays the basic role in the establishment of analytical planetary theories. Ancient astronomers: Pierce [1], Le Verrier [2], Boquet [3], Newcomb [16] expanded Δ^{-1} to degrees 6, 7, 8, 6 respectively, w.r.t. eccentricity, and mutual inclinations and nodes. We proceed two steps forward:

- We refer the inclinations and the nodes to a common fixed plane, say the ecliptic, which involve more complex computations, but lead to convenient results.
- By the present analysis we can calculate Δ^{-3} , Δ^{-5} , Δ^{-7} , ... by putting $s = 3, 5, 7, \dots$ in the development, which is essential, for the construction of planetary theories of order > 1 w.r.t. planetary masses.

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2. Methods and results

We can make the calculation of $\cos(f)$ and $\sin(f)$ to any order of e , by using the following formulas, for given s_m, i_m and j_m :

$$\begin{aligned} \cos(f) &\approx -e + \frac{2(1 - e^2)}{e} \cdot \sum_{s=1}^{s_m} J_s(se) \cos(sM) = \\ &-e + \frac{2(1 - e^2)}{e} \sum_{s=1}^{s_m} \sum_{i=0}^{i_m} \frac{(-1)^i \left(\frac{se}{2}\right)^{s+2i}}{i! (s+i)!} \cos(sM). \\ \sin(f) &\approx 2\sqrt{1 - e^2} \sum_{s=1}^{s_m} \frac{1}{s} \frac{d}{de} J_s(se) \sin(sM) \approx \\ &\left[-2 \sum_{j=0}^{j_m} \frac{(2j)! e^{2j}}{2^{2j} (j!)^2 (2j - 1)} \right] \cdot \sum_{s=1}^{s_m} \frac{1}{s} \left[\frac{s}{2} \int_{i=0}^{i_m} \frac{(-1)^i (s + 2i) \left(\frac{se}{2}\right)^{s+2i-1}}{i! (s+i)!} \right] \sin(sM). \\ \sin f &\approx - \sum_{j=0}^{j_m} \int_{s=1}^{s_m} \int_{i=0}^{i_m} \frac{(2j)! e^{2j}}{2^{2j} (j!)^2 (2j - 1)} \frac{(-1)^i (s + 2i) \left(\frac{se}{2}\right)^{s+2i-1}}{i! (s+i)!} \sin(sM). \end{aligned}$$

We have the following results concerning the calculation of $\cos(f)$ and $\sin(f)$ to e^6 obtained by taking $i_m = j_m = s_m = 6$.

$$\begin{aligned} \cos(f) &= \cos(M) + e[\cos(2M) - 1] + 9e^2 \frac{\cos(3M) - \cos(M)}{8} + \\ &4e^3 \frac{\cos(4M) - \cos(2M)}{3} + 25e^4 \frac{25 \cos(5M) - 27 \cos(3M) + 2 \cos(M)}{384} + \\ &3e^5 \frac{27 \cos(6M) - 32 \cos(4M) + 5 \cos(2M)}{40} + \\ &49e^6 \frac{[2401 \cos(7M) - 3125 \cos(5M) + 729 \cos(3M) - 5 \cos(M)]}{46080}. \\ \sin(f) &= \sin(M) + e \sin(2M) + e^2 \frac{9 \sin(3M) - 7 \sin(M)}{8} + \\ &e^3 \frac{8 \sin(4M) - 7 \sin(2M)}{6} + \\ &e^4 \frac{625 \sin(5M) - 621 \sin(3M) + 34 \sin(M)}{384} + \\ &e^5 \frac{243 \sin(6M) - 272 \sin(4M) + 40 \sin(2M)}{120} + \\ &e^6 \frac{117649 \sin(7M) - 146875 \sin(5M) + 33129 \sin(3M) - 1355 \sin(M)}{46080}. \end{aligned}$$

We have also the expansion series of r/a :

$$\frac{r}{a} \approx 1 + \frac{e^2}{2} - 2e \sum_{s=1}^{s_m} \frac{1}{s^2} \frac{d}{de} J_s(se) \cos(sM) \approx$$

$$1 - e \cos(M) + e^2 \frac{1 - \cos(2M)}{2} + e^3 \frac{3 \cos(M) - \cos(3M)}{8} + e^4 \frac{\cos(2M) - \cos(4M)}{3} - e^5 \frac{5[25 \cos(5M) + -27 \cos(3M) + 2 \cos(M)]}{384} + e^6 \frac{[27 \cos(6M) - 32 \cos(4M) + 5 \cos(2M)]}{80}$$

Then we deduce the expansion of

$$\varepsilon \equiv \frac{r}{a} - 1 \approx \frac{e^2}{2} - 2e \sum_{s=1}^{s_m} \frac{1}{s^2} \frac{d}{de} J_s(se) \cos(sM) \approx e \cos(M) + e^2 \frac{1 - \cos(2M)}{2} + e^3 \frac{3[\cos(M) - \cos(3M)]}{8} + e^4 \frac{[\cos(2M) - \cos(4M)]}{3} + e^5 \frac{5[25 \cos(5M) - 27 \cos(3M) + 2 \cos(M)]}{384} + e^6 \frac{27 \cos(6M) - 32 \cos(4M) + 5 \cos(2M)}{80}.$$

Hence we obtain the expansion of $\varepsilon' \equiv \frac{r'}{a'} - 1$:

$$\varepsilon' \approx -e' \cos(M') + e'^2 \frac{1 - \cos(2M')}{2} + e'^3 \frac{3[\cos(M') - \cos(3M')]}{8} + e'^4 \frac{\cos(2M') - \cos(4M')}{3} - e'^5 \frac{5[25 \cos(5M') - 27 \cos(3M') + 2 \cos(M')]}{384} - e'^6 \frac{27 \cos(6.M') - 32 \cos(4.M') + 5 \cos(2M')}{80}$$

Then we deduce

$$\cos(\omega + f) = \cos(\omega) \cos(f) - \sin(\omega) \sin(f)$$

and

$$\sin(\omega + f) = \sin(\omega) \cos(f) + \cos(\omega) \sin(f).$$

After calculation we obtain:

$$\begin{aligned} \cos(\omega + f) &= \cos(M + \omega) + e [\cos(2M + \omega) - \cos(\omega)] + e^2 \frac{9 \cos(3M + \omega) - 8 \cos(M + \omega) - \cos(M - \omega)}{8} + e^3 \frac{16 \cos(4M + \omega) - 15 \cos(2M + \omega) - \cos(2M - \omega)}{12} + e^4 \frac{625 \cos(5M + \omega) - 648 \cos(3M + \omega) - 27 \cos(3M - \omega) + 42 \cos(M + \omega) + 8 \cos(M - \omega)}{384} + \frac{e^5}{240} \left[\frac{486 \cos(6M + \omega) - 560 \cos(4M + \omega) - 16 \cos(4M - \omega) + 85 \cos(2M + \omega) + 5 \cos(2M - \omega)}{85 \cos(2M + \omega) + 5 \cos(2M - \omega)} \right] + \end{aligned}$$

$$\frac{e^6}{46080} \left[\begin{array}{l} 117649 \cos(7M + \omega) - 150000 \cos(5M + \omega) - 3125 \cos(5M - \omega) + \\ 34425 \cos(3M + \omega) + 1296 \cos(3M - \omega) - 800 \cos(M + \omega) + 555 \cos(M - \omega) \end{array} \right].$$

Then we have:

$$\begin{aligned} \sin(\omega + f) &= \sin(M + \omega) + e [\sin(2M + \omega) - \sin(\omega)] + \\ &e^2 \frac{9 \sin(3M + \omega) - 8 \sin(M + \omega) + \sin(M - \omega)}{8} + \\ &e^3 \frac{16 \sin(4M + \omega) - 15 \sin(2M + \omega) + \sin(2M - \omega)}{12} + \\ &e^4 \frac{625 \sin(5M + \omega) - 648 \sin(3M + \omega) + 27 \sin(3M - \omega) + 42 \sin(M + \omega) - 8 \sin(M - \omega)}{384} + \\ &e^5 \frac{486 \sin(6M + \omega) - 560 \sin(4M + \omega) + 16 \sin(4M - \omega) + 85 \sin(2M + \omega) - 5 \sin(2M - \omega)}{240} + \\ &e^6 \left[\frac{117649 \sin(7M + \omega) - 150000 \sin(5M + \omega) + 3125 \sin(5M - \omega)}{46080} + \right. \\ &\left. + \frac{34425 \sin(3M + \omega) - 1296 \sin(3M - \omega) - 800 \sin(M + \omega) - 555 \sin(M - \omega)}{46080} \right]. \end{aligned}$$

We have, on the other hand,

$$\begin{aligned} \frac{x}{r} &= \cos(\Omega) \cos(\omega + f) - \sin(\Omega) \sin(\omega + f) \cos(i) \\ \frac{y}{r} &= \sin(\Omega) \cos(\omega + f) + \cos(\Omega) \sin(\omega + f) \cos(i) \\ \frac{z}{r} &= \sin(\omega + f) \sin i \end{aligned}$$

and

$$\begin{aligned} \frac{x_1}{r_1} &= \cos(\Omega_1) \cos(\omega_1 + f_1) - \sin(\Omega_1) \sin(\omega_1 + f_1) \cos(i) \\ \frac{y_1}{r_1} &= \sin(\Omega_1) \cos(\omega_1 + f_1) + \cos(\Omega_1) \sin(\omega_1 + f_1) \cos(i) \\ \frac{z_1}{r_1} &= \sin(\omega_1 + f_1) \sin i \end{aligned}$$

which gives:

$$\cos(\psi) = \frac{xx_1 + yy_1 + zz_1}{rr_1}.$$

We have also:

$$\theta = \varpi + f = \Omega + \omega + f, \quad \theta_1 = \varpi_1 + f_1 = \Omega_1 + \omega_1 + f_1,$$

$$\theta - \theta_1 = (\Omega - \Omega_1) + [(\omega + f) - (\omega_1 + f_1)]$$

$$\begin{aligned} \cos(\theta - \theta_1) &= \cos[\Omega + \omega + f] \cos[\Omega_1 + \omega_1 + f_1] + \sin[\Omega + \omega + f] \sin[\Omega_1 + \omega_1 + f_1] = \\ &= [\cos \Omega \cos(\omega + f) - \sin(\Omega) \sin(\omega + f)][\cos(\Omega_1) \cos(\omega_1 + f_1) - \sin(\Omega_1) \sin(\omega_1 + f_1)] + \\ &[\sin(\Omega) \cos(\omega + f) + \cos(\Omega) \sin(\omega + f)][\sin(\Omega_1) \cos(\omega_1 + f_1) + \cos(\Omega_1) \sin(\omega_1 + f_1)] \end{aligned}$$

$$\begin{aligned}
 &= \cos(\Omega - \Omega_1)[\cos(\omega + f) \cos(\omega_1 + f_1) + \sin(\omega + f) \sin(\omega_1 + f_1)] + \\
 &\quad - \sin(\Omega - \Omega_1)[\sin(\omega + f) \cos(\omega_1 + f_1) - \cos(\omega + f) \sin(\omega_1 + f_1)] \\
 &\quad \sin(I/2) = \gamma, \quad \cos(I) = 1 - 2 \sin^2(I/2) = 1 - 2\gamma^2, \\
 \sin(I) &= 2\gamma\sqrt{1 - \gamma^2} \approx 2 \left(1 - \frac{\gamma^2}{2} - \frac{\gamma^4}{8} - \frac{\gamma^6}{16} - \frac{5\gamma^8}{128} \right) = 2 - \gamma^3 - \frac{\gamma^5}{4} - \frac{\gamma^7}{8} - \frac{5\gamma^9}{64} \\
 &\quad \approx 2 - \gamma^3 - \frac{\gamma^5}{4} - \frac{\gamma^7}{8}
 \end{aligned}$$

Let us use the longitudes $\lambda = M + \varpi$ and $\lambda_1 = M_1 + \varpi_1$, by replacing, in the expressions of $\cos(\psi)$ and $\cos(\theta - \theta_1)$, ω and M , ω_1 and M_1 respectively by:

$$\omega = \varpi - \Omega, \quad M = \lambda - \varpi, \quad \omega_1 = \varpi_1 - \Omega_1, \quad M_1 = \lambda_1 - \varpi_1.$$

The expression of $\cos(\theta - \theta_1)$ can be obtained from the expression of $\cos(\psi)$ by setting $\psi' = 0$. Then we can deduce the expression of

$$\Psi = \cos(\psi) - \cos(\theta - \theta_1).$$

We obtain:

$$\begin{aligned}
 \Psi &= -2e\gamma\gamma' \cos(2\lambda - \varpi + \lambda' - \varpi') + e\gamma'^2 \cos(2\lambda - \varpi + \lambda' - \varpi') \\
 &\quad + e\gamma^2 \cos[2\lambda - \varpi + \lambda' - 2\varpi] + 2e\gamma\gamma' \cos(2\lambda - \varpi - \lambda' + \varpi') \\
 &\quad - e(\gamma^2 + \gamma'^2) \cos(2\lambda - \varpi - \lambda') + 2e\gamma\gamma' \cos(\lambda + \varpi' - 2\lambda' + \varpi' - \varpi) \\
 &\quad - e'(\gamma^2 + \gamma'^2) \cos(\lambda + \varpi' - 2\lambda') + 2e'\gamma\gamma' \cos(\lambda + \varpi' - \varpi' - \varpi) \\
 &\quad - e'\gamma'^2 \cos(\lambda + \varpi' - 2') - e'\gamma^2 \cos(\lambda + \varpi' - 2) \\
 &\quad - 2e'\gamma\gamma' \cos(\lambda - \varpi' + 2\lambda' - \varpi' - \varpi) \\
 &\quad + e'\gamma'^2 \cos(\lambda - \varpi' + 2\lambda' - 2') + e'\gamma^2 \cos(\lambda - \varpi' + 2\lambda' - 2) \\
 &\quad - 2e'\gamma\gamma' \cos(\lambda - \varpi' + \varpi' - \varpi) + e'(\gamma^2 + \gamma'^2) \cos(\lambda - \varpi') \\
 &\quad - 2\gamma\gamma' \cos(\lambda + \lambda' - \varpi' - \varpi) \\
 &\quad + \gamma'^2 \cos(\lambda + \lambda' - 2') + \gamma^2 \cos(\lambda + \lambda' - 2) \\
 &\quad + 2\gamma\gamma' \cos(\lambda - \lambda' + \varpi' - \varpi) - (\gamma^2 + \gamma'^2) \cos(\lambda - \lambda') \\
 &\quad + e\gamma^2 \cos(\varpi - \lambda') - e\gamma^2 \cos(\varpi + \lambda' - 2) \\
 &\quad + 2e\gamma\gamma' \cos(\varpi + \lambda' - \varpi') - 2e\gamma\gamma' \cos(\varpi - \lambda' + \varpi' - \varpi) \\
 &\quad - e\gamma'^2 \cos(\varpi + \lambda' - 2') + e\gamma'^2 \cos(\varpi - \lambda')
 \end{aligned}$$

We note that Ψ is of second order in γ , γ' , e and e' . The terms of the second order are independent of e and e' .

On the other hand, we have

$$\cos \theta = \cos(\Omega + \omega + f) = \cos(\varpi + f) = \cos(\varpi) \cos(f) - \sin(\varpi) \sin(f),$$

$$\sin \theta = \sin(\Omega + \omega + f) = \sin(\varpi + f) = \sin(\varpi) \cos(f) + \cos(\varpi) \sin(f).$$

Then we can deduce

$$e^{i\theta} = \cos(\theta) + i \sin(\theta),$$

$$e^{ij\theta} = (\cos(\theta) + i \sin(\theta))^j = \cos(j\theta) + i \sin(j\theta),$$

so we obtain $\cos(j\theta)$ and $\sin(j\theta)$ as series of ϖ , λ and M . We notice that it is possible to calculate f as a function of M . For this aim, we use the differential relation $dM = n dt$, with $n = \sqrt{\frac{\mu}{a^3}}$, $M = E - e \sin(E)$,

$$dM = (1 - e \cos(E)) dE$$

$$df = \frac{\sqrt{\mu a(1 - e^2)}}{r^2} dt = \frac{\sqrt{\mu a(1 - e^2)}}{a(1 - e \cos(E))^2} dt,$$

with $r = a(1 - e \cos(E))$,

$$df = \sqrt{\frac{\mu}{a^3}} \frac{\sqrt{1 - e^2}}{(1 - e \cos(E))} dt = \frac{\sqrt{1 - e^2} dM}{(1 - e \cos E)^2} = \sqrt{1 - e^2} \left(\frac{dE}{dM} \right)^2 dM.$$

We have also

$$dM = \frac{nr^2 df}{\sqrt{\mu a(1 - e^2)}} = (1 - e^2)^{3/2} \frac{df}{(1 + e \cos(f))^2}.$$

We develop the last member as a series before integrating term by term, the constant of integration being equal to zero:

$$M = f - 2e \sin(f) + \frac{3e^2 \sin(2f)}{4} - \frac{e^3 \sin(3f)}{3} + e^4 \left[\frac{5 \sin(4f)}{32} + \frac{\sin(4f)}{8} \right] +$$

$$-e^5 \left[\frac{3 \sin(5f)}{40} + \frac{\sin(3f)}{8} \right] + e^6 \left[\frac{7 \sin(6f)}{192} + \frac{3 \sin(4f)}{32} + \frac{3 \sin(2f)}{64} \right] + \dots$$

We can deduce: $f = M + e\Phi(f)$, with

$$\Phi(f) \equiv \frac{f - M}{e} = 2 \sin(f) + \frac{3e \sin(2f)}{4} + \frac{e^2 \sin(3f)}{3} - e^3 \left[\frac{5 \sin(4f)}{32} + \frac{\sin(4f)}{8} \right]$$

$$+ e^4 \left[\frac{3 \sin(5f)}{40} + \frac{\sin(3f)}{8} \right] - e^5 \left[\frac{7 \sin(6f)}{192} + \frac{3 \sin(4f)}{32} + \frac{3 \sin(2f)}{64} \right]$$

$$+ e^6 \left[\frac{\sin(7f)}{56} + \frac{\sin(5f)}{16} + \frac{\sin(3f)}{16} \right] + \dots$$

So we can write.

$$f = M + \sum_{j=1}^{\infty} \frac{d^{j-1}}{dM^{j-1}} [\Phi(M)]^j$$

$$f = M + 2e \sin(M) + \frac{5e^2 \sin(2M)}{4} + e^3 \left[\frac{13 \sin(3M)}{12} - \frac{\sin(M)}{4} \right] +$$

$$+e^4 \left[\frac{103 \sin(4M)}{96} - \frac{11 \sin(2M)}{24} \right] + e^5 \left[\frac{1097 \sin(5M)}{960} - \frac{43 \sin(3M)}{64} + \frac{5 \sin(M)}{96} \right] \\ + e^6 \left[\frac{1223 \sin(6M)}{960} - \frac{451 \sin(4M)}{480} + \frac{17 \sin(2M)}{192} \right]$$

We notice that $(f - M)$ is of order 1 at least with respect to e . We have:

$$j(f - M) = 2je \sin(M) + \frac{5je^2 \sin(2M)}{4} + je^3 \left[\frac{13 \sin(3M)}{12} - \frac{\sin(M)}{4} \right] \\ + e^4 \left[\frac{103 \sin(4M)}{96} - \frac{11 \sin(2M)}{24} \right] + je^5 \left[\frac{1097 \sin(5M)}{960} - \frac{43 \sin(3M)}{64} + \frac{5 \sin(M)}{96} \right] \\ + je^6 \left[\frac{1223 \sin(6M)}{960} - \frac{451 \sin(4M)}{480} + \frac{17 \sin(2M)}{192} \right]$$

$$\cos((j(f - M))) = \\ = \sum_{i=0}^{\infty} \frac{(-1)^i [j(f - M)]^{2i}}{(2i)!} \approx 1 - \frac{[j(f - M)]^2}{2} + \frac{[j(f - M)]^4}{24} + \frac{[j(f - M)]^6}{720}$$

to order 6.

$$\approx 1 - \frac{[j(f - M)]^2}{2}$$

to order 3.

$$\sin((j(f - M))) = \\ = \sum_{i=0}^{\infty} \frac{(-1)^i [j(f - M)]^{2i+1}}{(2i + 1)!} \approx 1 - j(f - M) - \frac{[j(f - M)]^3}{6} + \frac{[j(f - M)]^5}{120}$$

to order 6.

$$\approx [1 - j(f - M)] - \frac{[j(f - M)]^3}{6}$$

to order 3.

We obtain to order 3:

$$\cos(j(f - M)) \approx 1 + e^2 j^2 [\cos(2M) - 1] + \frac{5}{4} e^3 j^2 [\cos(3M) - \cos(M)] \\ \sin[j(f - M)] \approx 2ej \sin(M) + \frac{5}{4} e^2 j \sin(2M) + \\ + \frac{1}{12} e^3 j [(4j^2 + 13 \sin(3M) - 3(4j^2 + 1 \sin(M))]$$

On the other hand, we can write, since: $\theta = \omega + \Omega + f$ and $M = \lambda - \varpi$:

$$\cos(j\theta) \equiv \cos[j(\omega + \Omega + f)] = \cos[j(\varpi + f)] = \cos[j\varpi + jM + j(f - M)] = \\ = \cos[j\lambda + j(f - M)]$$

$$\begin{aligned}
&= \cos(j\lambda) \cos[j(f - M)] - \sin(j\lambda) \sin[j(f - M)] = \\
\sin(j\theta) &\equiv \sin[j(\omega + \Omega + f)] = \sin[j(\varpi + f)] = \sin[j\varpi + jM + j(f - M)] = \\
&\sin[j\lambda + j(f - M)] = \sin(j\lambda) \cos[j(f - M)] + \cos(j\lambda) \sin[j(f - M)],
\end{aligned}$$

We find, after calculation:

$$\begin{aligned}
\cos(j\theta) &\approx \cos(j\lambda) + e [j \cos[(j + 1)\lambda - \varpi] - \cos[(j - 1)\lambda + \varpi]] \\
&+ \frac{e^2 j}{8} \{ (4j + 5) \cos[(j + 2)\lambda - 2\varpi] + (4j - 5) \cos[(j - 2)\lambda + 2\varpi] - 8j \cos(j\lambda) \} \\
&+ \frac{e^3 j}{24} \{ (4j^2 + 15j + 13) \cos[(j + 3)\lambda - 3\varpi] - 3(4j^2 + 5j + 1) \cos[(j + 1)\lambda - \varpi] \\
&\quad + 3(4j^2 - 5j + 1) \cos[(j - 1)\lambda + \varpi] - (4j^2 - 15j + 13) \cos[(j - 1)\lambda + \varpi] \} \\
\sin(j\theta) &\approx \sin(j\lambda) + ej \{ \sin[(j + 1)\lambda - \varpi] - \sin[(j - 1)\lambda + \varpi] \} \\
&+ \frac{e^2 j}{8} \{ (4j + 5) \sin[(j + 2)\lambda - 2\varpi] + (4j - 5) \sin[(j - 2)\lambda + 2\varpi] - 8j \sin(j\lambda) \} \\
&+ \frac{e^3 j}{24} \{ (4j^2 + 15j + 13) \sin[(j + 3)\lambda - 3\varpi] - 3(4j^2 + 5j + 1) \sin[(j + 1)\lambda - \varpi] \\
&\quad + 3(4j^2 - 5j + 1) \sin[(j - 1)\lambda + \varpi] - (4j^2 - 15j + 13) \sin[(j - 1)\lambda + \varpi] \}
\end{aligned}$$

So we can calculate:

$$\begin{aligned}
\cos[j(\theta - \theta')] &= \cos(j\theta) \cos(j\theta') + \sin(j\theta) \sin(j\theta') \\
&= \frac{1}{24} \{ e^3 j(4j^2 + 15j + 13) \cos[(j + 3)\lambda - 3\varpi - j\lambda'] \\
&\quad + 3e^2 e' j^2 (4j + 5) \cos[(j + 2)\lambda - (j + 1)\lambda' + \varpi' - 2\varpi] \\
&\quad - 3e^2 e' j^2 (4j + 5) \cos[(j + 2)\lambda - (j - 1)\lambda' - \varpi' - 2\varpi] \\
&\quad + 3e^2 j(4j + 5) \cos[(j + 2)\lambda - j\lambda' - 2\varpi] \\
&\quad + 3ee'^2 j^2 (4j + 5) \cos[(j + 1)\lambda - (j + 2)\lambda' + 2\varpi' - \varpi] \\
&\quad + 24ee' j^2 \cos[(j + 1)(\lambda - \lambda') + \varpi' - \varpi] \\
&\quad - 24ee' j^2 \cos[(j + 1)\lambda - (j - 1)\lambda' - \varpi' - \varpi] \\
&\quad + 3ee'^2 j^2 (4j - 5) \cos[(j + 1)\lambda - (j - 2)\lambda' - 2\varpi' - \varpi] \\
&\quad + 3ej[-4(2e'^2 + e^2)j^2 - e^2(5j + 1) + 8] \cos[(j + 1)\lambda - j\lambda' - \varpi] \\
&\quad + e'^3 j(4j^2 + 15j + 13) \cos[j\lambda - (j + 3)\lambda' + 3\varpi'] \\
&\quad + 3e'^2 j(4j + 5) \cos[j\lambda - (j + 2)\lambda' + 2\varpi'] \\
&\quad - 3e' j[4(2e^2 + e'^2)j^2 + e'^2(5j + 1) - 8] \cos[j\lambda - (j + 1)\lambda' + \varpi'] \\
&\quad + 3e' j[4(2e^2 + e'^2)j^2 - e'^2(5j - 1) - 8] \cos[j\lambda - (j - 1)\lambda' - \varpi'] \\
&\quad + 3e'^2 j(4j - 5) \cos[j\lambda - (j - 2)\lambda' - 2\varpi']
\end{aligned}$$

$$\begin{aligned}
 & -e'^3 j(4j^2 - 15j + 13) \cos[j\lambda - (j - 3)\lambda' - 3\varpi'] \\
 & -24[(e^2 + e'^2)j^2 - 1] \cos[j(\lambda - \lambda')] \\
 & -3ee'^2 j^2(4j + 5) \cos[(j - 1)\lambda - (j + 2)\lambda' + 2\varpi' + \varpi] \\
 & -24ee' j^2 \cos[(j - 1)\lambda - (j + 1)\lambda' + \varpi' + \varpi] \\
 & +24ee' j^2 \cos[(j - 1)\lambda - (j - 1)\lambda' - \varpi' + \varpi] \\
 & -3ee'^2 j^2(4j - 5) \cos[(j - 1)\lambda - (j - 2)\lambda' - 2\varpi' + \varpi] \\
 & +3ej[(e^2 + 2e'^2)j^2 - e^2(5j - 1) - 8] \cos[(j - 1)\lambda - j\lambda' + \varpi] \\
 & +3e^2 e' j^2(4j - 5) \cos[(j - 2)\lambda - (j + 1)\lambda' + 2\varpi + \varpi'] \\
 & -3e^2 e' j^2(4j - 5) \cos[(j - 2)\lambda - (j - 1)\lambda' + 2\varpi - \varpi'] \\
 & +3e^2 j(4j - 5) \cos[(j - 2)\lambda - j\lambda' + 2\varpi] \\
 & -e^3 j(4j^2 - 5j + 13) \cos[(j - 3)\lambda - j\lambda' + 3\varpi] \}
 \end{aligned}$$

In the next step, we calculate $\left(\frac{r}{a} \frac{r'}{a'} \Psi\right)$. We obtain to order 3:

$$\begin{aligned}
 \left(\frac{r}{a} \frac{r'}{a'} \Psi\right) &= \frac{rr'\Psi}{aa'} = -e \cos(2\lambda + \lambda' - \varpi - \varpi') + (e'^2/2) \cos(2\lambda + \lambda' - \varpi - 2\Omega') \\
 & + (e^2/2) \cos(2\lambda + \lambda' - \varpi - 2\Omega + e' \cos(2\lambda - \lambda' - \varpi - \Omega + \Omega')) \\
 & - [(e^2 + e'^2)/2] \cos(2\lambda - \lambda' - \varpi) + e'' \cos(\lambda - 2\lambda' + \varpi' - \Omega + \Omega') \\
 & - [e'(e^2 + e'^2)/2] \cos(\lambda - 2\lambda' + \varpi') + 3e'' \cos(\lambda + \varpi' - \Omega - \Omega') \\
 & - 3(e'^2/2) \cos(\lambda + \varpi' - 2\Omega') - 3[e'^2/2] \cos(\lambda + \varpi' - 2\Omega) \\
 & - e'' \cos(\lambda + 2\lambda' - \varpi' - \Omega - \Omega') \\
 & + [e''^2/2] \cos(\lambda + 2\lambda' - \varpi' - 2\Omega') + [e'^2/2] \cos(\lambda + 2\lambda' - \varpi' - 2\Omega) \\
 & - 3e'' \cos(\lambda - \varpi' - \Omega + \Omega') \\
 & + 3[e'(e^2 + e'^2)/2] \cos(\lambda - \varpi') - 2' \cos(\lambda + \lambda' - \Omega - \Omega') \\
 & + e'^2 \cos(\lambda + \lambda' - 2\Omega') + e^2 \cos(\lambda + \lambda' - 2\Omega) + 2' \cos(\lambda - \lambda' - \Omega + \Omega') \\
 & - (e^2 + e'^2) \cos(\lambda - \lambda') + [3e(e^2 + e'^2)/2] \cos(\lambda' - \varpi) \\
 & - [e^2/2] \cos(\lambda' + \varpi - 2\Omega) \\
 & + 3e' \cos(\lambda' + \varpi - \Omega - \Omega') - 3e' \cos(\lambda' - \varpi + \Omega - \Omega') \\
 & - [3e'^2/2] \cos(\lambda' + \varpi - 2\Omega').
 \end{aligned}$$

We note that $\left(\frac{r}{a} \frac{r'}{a'} \Psi\right)$ is of the second order in λ' and λ at least. Consequently:

$$\left(\frac{r}{a} \frac{r'}{a'} \Psi\right)^2$$

is of order 4 at least, and to order 3 we have:

$$\left(\frac{r r'}{a a'} \Psi\right)^2 \approx 0.$$

We have:

$$\begin{aligned} \Delta^{-s} &= \left(r^2 + r'^2 - 2rr' \cos \psi\right)^{-s/2} \\ &\equiv \left[r^2 + r'^2 - 2rr' \Psi + \cos(\theta - \theta')\right]^{-s/2} \\ &\equiv \left\{[r^2 + r'^2 - 2rr' \cos(\theta - \theta')] - 2rr' \Psi\right\}^{-s/2} \\ &\equiv [\Delta_0^2 - 2rr' \Psi]^{-s/2} = \Delta_0^{-s} \left[1 - \frac{2rr' \Psi}{\Delta_0^2}\right]^{-s/2} \\ &= \Delta_0^{-s} \sum_{k=0}^{\infty} \frac{(rr' \Psi)^k}{k! \Delta_0^{2k}} \prod_{l=0}^{k-1} (s + 2l) = \sum_{k=0}^{\infty} \prod_{l=0}^{k-1} (s + 2l) \frac{(rr' \Psi)^k}{k!} \Delta_0^{-(2k+s)}, \end{aligned}$$

We use the formula:

$$(1 - u)^{-\frac{s}{2}} = \int_{k=0}^{\infty} \frac{u^k}{2^k k!} \int_{l=0}^{k-1} (s + 2l),$$

with

$$u = \frac{2rr' \Psi}{\Delta_0^2}, \quad \left(\frac{u}{2}\right)^k = (rr' \Psi)^k,$$

where

$$\prod_{l=0}^{k-1} (s + 2l) = 1,$$

so we have:

$$(1 - u)^{-s/2} = 1 + \frac{su}{2(1!)} + \frac{s(s+2)}{(2!)} \left(\frac{u}{2}\right)^2 + \dots,$$

with

$$\begin{aligned} \Delta_0^2 &= \left[r^2 + r'^2 - 2rr' \cos(\theta - \theta')\right], \\ \Delta_0^{-(2k+s)} &= \left[r^2 + r'^2 - 2rr' \cos(\theta - \theta')\right]^{-(2k+s)/2}, \\ \Psi &= \cos \psi - \cos(\theta - \theta') \end{aligned}$$

and

$$(\theta - \theta') = (\varpi + f) - (\varpi' + f') = (\varpi - \varpi') + (f - f')$$

whence after some algebra:

$$\Delta^{-s} = \left(\frac{1}{\Delta}\right)^s = \sum_{k=0}^{\infty} \prod_{l=0}^{k-1} (s + 2l) \frac{(rr' \Psi)^k}{k!} \left(\frac{1}{\Delta_0}\right)^{(2k+s)},$$

with

$$\prod_{l=0}^{-1} (s + 2l) = 1,$$

where

$$\frac{1}{\Delta_0} = \left[r^2 + r'^2 - 2rr' \cos(\theta - \theta') \right]^{-1/2}$$

We remark that Ψ is of order 2 at least with respect to γ and γ' . Furthermore the Ψ terms of order 2 depend on the inclinations not on the eccentricities, so $(rr'\Psi)$ is of order 2 w.r.t. γ and γ' .

Hence to order 3, $k = 0$ and 1:

$$\begin{aligned} \left(\frac{1}{\Delta}\right)^s &= \sum_{k=0}^1 \frac{(rr'\Psi)^k}{k!} \left(\frac{1}{\Delta_0}\right)^{(2k+s)} \prod_{l=0}^{k-1} (s + 2l), \\ \left(\frac{1}{\Delta}\right)^s &= \frac{1}{\Delta_0^s} + (rr'\Psi) \left(\frac{1}{\Delta_0}\right)^{(2+s)} \quad s = \frac{1}{\Delta_0^s} \left[1 + s (rr'\Psi) \left(\frac{1}{\Delta_0}\right)^2 \right] \end{aligned}$$

Therefore, we have:

$$\left(\frac{1}{\Delta_0}\right)^{2k+s} = \left[r^2 + r'^2 - 2rr' \cos(\theta - \theta') \right]^{-(2k+s)/2}.$$

Let

$$\begin{aligned} \frac{1}{\rho_0} &= \left[a^2 + a'^2 - 2aa' \cos(\theta - \theta') \right]^{-1/2} \\ &= a'^{-1} \left[1 + \alpha^2 - 2\alpha \cos(\theta - \theta') \right]^{-1/2}, \quad \alpha = \frac{a}{a'}, \end{aligned}$$

so we have:

$$\left(\frac{1}{\rho_0}\right)^{2k+s} = a'^{-(2k+s)} \left[1 + \alpha^2 - 2\alpha \cos(\theta - \theta') \right]^{-(k+s/2)}$$

Applying Taylor's series expansion in ρ_0 , we may write

$$\frac{1}{\Delta_0} = \sum_{i=0}^{\infty} \frac{1}{i!} \left[(r - a) \frac{\partial}{\partial a} + (r' - a') \frac{\partial}{\partial a'} \right]^i \left(\frac{1}{\rho_0}\right)$$

We have also:

$$\left(\frac{1}{\Delta_0}\right)^{(2k+s)} = \sum_{i=0}^{\infty} \frac{1}{i!} \left[(r - a) \frac{\partial}{\partial a} + (r' - a') \frac{\partial}{\partial a'} \right]^i \left(\frac{1}{\rho_0}\right)^{(2k+s)}$$

Then we deduce:

$$\begin{aligned} \left(\frac{1}{\Delta}\right)^s &= \\ \sum_{k=0}^{\infty} \sum_{i=0}^{\infty} \frac{(rr'\Psi)^k}{k!} \frac{\prod_{l=0}^{k-1} (2.l + s)}{i!} &\left\{ \left[(r - a) \frac{\partial}{\partial a} + (r' - a') \frac{\partial}{\partial a'} \right]^i \left(\frac{1}{\rho_0}\right)^{(2k+s)} \right\} \end{aligned}$$

$$= \sum_{i=0}^{\infty} \sum_{k=0}^{\infty} \frac{(rr'\Psi)^k}{k!i!} \frac{\prod_{l=0}^{k-1} (2l+s)}{i!} \left[(r-a) \frac{\partial}{\partial a} + (r'-a') \frac{\partial}{\partial a'} \right]^i \left[a^2 + a'^2 - 2aa' \cos(\theta - \theta') \right]^{-(2k+s)/2}$$

But, we have:

$$\left(\frac{1}{\rho_0} \right)^v = \left[a^2 + a'^2 - 2aa' \cos(\theta - \theta') \right]^{-v/2} = a'^{-v} \left[1 + \alpha^2 - 2\alpha \cos(\theta - \theta') \right]^{-v/2} = \frac{a'^{-v}}{2} \sum_{j=-\infty}^{\infty} b_{v/2}^{(j)} \cos[j(\theta - \theta')] = \sum_{j=-\infty}^{\infty} B_{v/2}^{(j)} \cos[j(\theta - \theta')],$$

with

$$B_{v/2}^{(j)} = \frac{a'^{-v}}{2} b_{v/2}^{(j)},$$

where $b_{v/2}^{(j)}$ is a function of $\alpha = a/a'$, for all j . We have in particular

$$B_{(2k+s)/2}^{(j)} = \frac{a'^{-(2k+s)}}{2} b_{(2k+s)/2}^{(j)}(\alpha)$$

for all j with:

$$\left(\frac{1}{\rho_0} \right)^{2k+s} = a'^{-(2k+s)} \left[1 + \alpha^2 - 2\alpha \cos(\theta - \theta') \right]^{-(2k+s)/2} = \frac{a'^{-(2k+s)}}{2} \sum_{j=-\infty}^{\infty} b_{(2k+s)/2}^{(j)}(\alpha) \cos[j(\theta - \theta')]$$

Then we deduce:

$$\left(\frac{1}{\Delta} \right)^s = \sum_{i=0}^{\infty} \sum_{k=0}^{\infty} \sum_{j=-\infty}^{\infty} \frac{(rr'\Psi)^k}{k!i!} \sum_{l=0}^{k-1} (2l+s) \cdot$$

$$\cos[j(\theta - \theta')] \left[(r-a) \frac{\partial}{\partial a} + (r'-a') \frac{\partial}{\partial a'} \right]^i B_{(2k+s)/2}^{(j)}$$

or

$$\left(\frac{1}{\Delta} \right)^s = \sum_{i=0}^{\infty} \sum_{k=0}^{\infty} \sum_{j=-\infty}^{\infty} \prod_{l=0}^{k-1} (2l+s) \frac{(rr'\Psi)^k}{k!i!} \cos[j(\theta - \theta')] \cdot \left[(r-a) \frac{\partial}{\partial a} + (r'-a') \frac{\partial}{\partial a'} \right]^i \left[\frac{a' - (2k+s)}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right]$$

But we can write, for any function $c(a, a')$:

$$\left[(r-a) \frac{\partial}{\partial a} + (r'-a') \frac{\partial}{\partial a'} \right]^i c(a, a')$$

$$= \sum_{p=0}^i C_i^p (r-a)^p (r'-a')^{i-p} \frac{\partial^p}{\partial a^p} \left[\frac{\partial^{i-p}}{\partial a'^{i-p}} c(a, a') \right],$$

where $C_i^p = i!/p!(i-p)!$. So we have:

$$\begin{aligned} \left(\frac{1}{\Delta}\right)^s &= \sum_{i=0}^{\infty} \sum_{k=0}^{\infty} \sum_{j=-\infty}^{\infty} \frac{\prod_{l=0}^{k-1} (2l+s) (rr'\Psi)^k}{k!i!} \cos [j(\theta - \theta')]. \\ &\sum_{p=0}^i C_i^p (r-a)^p (r'-a')^{i-p} \frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+s)}}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right], \\ \left(\frac{1}{\Delta}\right)^s &= \sum_{i=0}^{\infty} \sum_{k=0}^{\infty} \sum_{j=-\infty}^{\infty} \frac{\prod_{l=0}^{k-1} (2l+s) (rr'\Psi)^k}{k!i!} \cos [j(\theta - \theta')] \sum_{p=0}^i C_i^p \varepsilon^p \varepsilon'^{i-p} a^p a'^{i-p} \\ &\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+s)}}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] \end{aligned}$$

with $\varepsilon = r/a - 1$, $\varepsilon' = r'/a' - 1$.

Finally, we obtain:

$$\begin{aligned} \left(\frac{1}{\Delta}\right)^s &= \sum_{j=-\infty}^{\infty} \sum_{k=0}^{\infty} \int_{i=0}^{\infty} \int_{p=0}^i a^p a'^{i-p} \frac{\prod_{l=0}^{k-1} (2l+s) \varepsilon^p \varepsilon'^{i-p} (rr'\Psi)^k \cos [j(\theta - \theta')]}{k!p!(i-p)!} \\ &\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+s)}}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right], \end{aligned}$$

where

$$\begin{aligned} &\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+s)}}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] = \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-2k+p+s}}{2} \frac{d^p}{d\alpha^p} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] \\ &= \sum_{l=0}^{i-p} C_{i-p}^l \prod_{m=1}^l (2k+s+p+m) \left[\frac{a'^{-(2k+p+s+l)}}{2} \right] \frac{\partial^{i-p-l}}{\partial a'^{i-p-l}} \left[\frac{d^p}{d\alpha^p} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right]. \end{aligned}$$

We remark that:

$$\begin{aligned} &\frac{\partial}{\partial a'} \left[\frac{d^p}{d\alpha^p} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] = -\frac{\alpha}{a'} \left[\frac{d^{p+1}}{d\alpha^{p+1}} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] \\ &\frac{\partial^2}{\partial a'^2} \left[\frac{d^p}{d\alpha^p} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] = \frac{1}{a'^2} \left[2\alpha \frac{d^{p+1}}{d\alpha^{p+1}} b_{\frac{2k+s}{2}}^{(j)}(\alpha) + \alpha^2 \frac{d^{p+1}}{d\alpha^{p+1}} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right], \\ &\frac{\partial^3}{\partial a'^3} \left[\frac{d^p}{d\alpha^p} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] = \\ &-\frac{1}{a'^3} \left[6\alpha \frac{d^{p+1}}{d\alpha^{p+1}} b_{\frac{2k+s}{2}}^{(j)}(\alpha) + 6\alpha^2 \frac{d^{p+2}}{d\alpha^{p+2}} b_{\frac{2k+s}{2}}^{(j)}(\alpha) + \alpha^3 \frac{d^{p+3}}{d\alpha^{p+3}} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right], \dots \end{aligned}$$

The expression of $(\frac{1}{\Delta})^s$ can be approximated by the truncated series:

$$g(s, i_m, j_m, k_m) = \sum_{i=0}^{i_m} \sum_{k=0}^{k_m} \sum_{j=-j_m}^{j_m} \sum_{p=0}^i a^p a'^{i-p} \frac{\prod_{l=0}^{k-1} (2l + s) \varepsilon^p \varepsilon'^{i-p} (rr'\Psi)^k}{k!p!(i-p)!} \cos [j(\theta - \theta')] ,$$

$$\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+s)}}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] .$$

In particular, we have for $s = 1$, since $\prod_{l=0}^{k-1} (2l + 1) = \frac{(2k)!}{(k!)2^k}$:

$$\frac{1}{\Delta} = \sum_{i=0}^{\infty} \sum_{k=0}^{\infty} \sum_{j=-\infty}^{\infty} \sum_{p=0}^i a^p a'^{i-p} \frac{(2k)! \varepsilon^p \varepsilon'^{i-p} (rr'\Psi/2)^k}{(k!)^2 p!(i-p)!} \cos [j(\theta - \theta')] .$$

$$\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+1)}}{2} b_{k+\frac{1}{2}}^{(j)}(\alpha) \right] ,$$

$$\frac{1}{\Delta} \cong g(1, i_m, j_m, k_m) =$$

$$\sum_{j=-j_m}^{j_m} \sum_{k=0}^{k_m} \sum_{i=0}^{i_m} \sum_{p=0}^i a^p a'^{i-p} \frac{(2k)! \varepsilon^p \varepsilon'^{i-p} (rr'\Psi/2)^k \cos [j(\theta - \theta')]}{k!p!(i-p)!}$$

$$\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+1)}}{2} b_{k+\frac{1}{2}}^{(j)}(\alpha) \right]$$

$$= \sum_{j=-j_m}^{j_m} \sum_{k=0}^{k_m} \sum_{i=0}^{i_m} \sum_{p=0}^i$$

$$\frac{(2k)! a^{p+k} a'^{i-p+k}}{(k!)^2 p!(i-p)!} \left(\frac{rr'\Psi}{2aa'} \right)^k \varepsilon^p \varepsilon'^{i-p} \frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \left[\frac{a'^{-(2k+1)}}{2} b_{k+\frac{1}{2}}^{(j)}(\alpha) \right] \cos [j(\theta - \theta')]$$

where:

$$\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} a'^{i-p} \left[\frac{a'^{-(2k+1)}}{2} b_{k+\frac{1}{2}}^{(j)}(\alpha) \right] = \frac{\partial^{i-p}}{\partial a'^{i-p}} \frac{a'^{-(2k+p+1)}}{2} \frac{d^p}{d\alpha^p} \left[b_{k+\frac{1}{2}}^{(j)}(\alpha) \right] = \dots$$

Consequently, to have an approximation of $(\frac{1}{\Delta})^s$ to order six, we must proceed as follows.

1. We calculate, to order 6, the expression of the product

$$r_c[k, j] = (rr'\Psi)^k \cos [j(\theta - \theta')] .$$

2. We calculate:

$$\begin{aligned}
 pr_j(s, i, k, p) &= \frac{\prod_{l=0}^{k-1} (2l + s)}{p!k!(i-p)!} a^p a'^{i-p} \varepsilon^p \varepsilon'^{i-p} \left[\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} B_{k+\frac{s}{2}}^{(j)} \right] \\
 &= \alpha^p \frac{\prod_{l=0}^{k-1} (2l + s)}{p!k!(i-p)!} \varepsilon^p \varepsilon'^{i-p} a'^i \left[\frac{\partial^{i-p}}{\partial a'^{i-p}} \left[a'^{-p} \frac{d^p}{d\alpha^p} \right] B_{k+\frac{s}{2}}^{(j)} \right]
 \end{aligned}$$

with $\varepsilon = \frac{r}{a} - 1$ and $\varepsilon' = \frac{r'}{a'} - 1$. We remark that $r_c[k, j]$ is of order $(2k)$ because $\cos [j(\theta - \theta')]$ contains a constant term (independent of parameters $e, e',$ and γ nad γ'). The product $\varepsilon^p \varepsilon'^{i-p}$ is of order i because ε and ε' are of order one with respect to e and e' respectively.

Consequently the product $pc_j[s, i, k, p] = r_c[k, j]pr_j(s, i, k, p)$ is of order $(2k + i)$.

To have the terms of order 2, we must take $0 \leq 2k + i \leq 2$, with $0 \leq p \leq i$. So the only terms of order ≤ 2 are

$$\begin{aligned}
 pc_j(s, 0, 0, 0), \quad pc_j(s, 0, 1, 0), \quad pc_j(s, 1, 0, 0), \quad pc_j(s, 1, 0, 1), \quad pc_j(s, 2, 0, 0), \\
 pc_j(s, 2, 0, 1) \quad \text{and} \quad pc_j(s, 2, 0, 2).
 \end{aligned}$$

3. We calculate

$$g(s, i_m, j_m, k_m) = \sum_{i=0}^{i_m} \sum_{k=0}^{k_m} \sum_{j=-j_m}^{j_m} \sum_{p=0}^i pc_j[s, i, k, p] = \sum_{j=-j_m}^{j_m} g_j(s, i_m, k_m),$$

with

$$\begin{aligned}
 g_j(s, i_m, k_m) &= \sum_{i=0}^{i_m} \sum_{k=0}^{k_m} \sum_{p=0}^i pc_j[s, i, k, p] \\
 &= \sum_{i=0}^{i_m} \sum_{k=0}^{k_m} \sum_{p=0}^i a'^i \alpha^p \frac{\prod_{l=0}^{k-1} (2l + s)}{p!k!(i-p)!} \varepsilon^p \varepsilon'^{i-p} \left[\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} B_{k+\frac{s}{2}}^{(j)} \right] (rr'\Psi)^k \cos [j(\theta - \theta')]
 \end{aligned}$$

where

$$\begin{aligned}
 pc_j[s, i, k, p] &= pr_j[s, i, k, p] r_c[k, j] \\
 &= a'^i \alpha^p \frac{\prod_{l=0}^{k-1} (2l + s)}{p!k!(i-p)!} \varepsilon^p \varepsilon'^{i-p} \left[\frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} B_{k+\frac{s}{2}}^{(j)} \right] (rr'\Psi)^k \cos [j(\theta - \theta')]
 \end{aligned}$$

It is sufficient to take $k_m = 3$, to obtain the series to order 6, since $r_c[k, j] = 0$, for $k > 3$.

[We will not use the expansion of Laplace coefficients $B(s, j)$ with respect to $\alpha = a/a'$. These coefficients will appear with their derivatives in the coefficients of the cosine terms as well as e and γ . When we require numerical

results we should assign the per limit of power of alpha, according to the number of decimals we desire.]

On the other hand, $i_m = 6, i = 0, 1, 2, 3, \dots, 6$. Hence the approximation of $(\frac{1}{\Delta})^s$ to order six is $g(s, 6, 3) = \sum_{j=-j_m}^{j_m} g_j(s, 6, 3)$, with

$$g_j(s, 6, 3) = \sum_{i=0}^6 \sum_{k=0}^3 \sum_{p=0}^i pr_j[s, i, k, p] r_c[k, j] = \sum_{i=0}^6 \sum_{k=0}^3 \sum_{p=0}^i pc_j[s, i, k, p],$$

where $pc_j[s, i, k, p] = pr_j[s, i, k, p] r_c[k, j]$

$$\begin{aligned} &= a^p a'^{i-p} \frac{\prod_{l=0}^{k-1} (2l + s) \varepsilon^p \varepsilon'^{i-p} (rr'\Psi)^k \cos [j(\theta - \theta')]}{k!p!(i-p)!} \frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \\ &\quad \left[\frac{a'^{-(2k+s)}}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] \\ &= \alpha^p a'^i \frac{\prod_{l=0}^{k-1} (2l + s) \varepsilon^p \varepsilon'^{i-p} (rr'\Psi)^k \cos [j(\theta - \theta')]}{k!p!(i-p)!} \frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a'^{i-p}} \\ &\quad \left[\frac{a'^{-(2k+s)}}{2} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] \\ &= \alpha^p a'^i \frac{\prod_{l=0}^{k-1} (2l + s) \varepsilon^p \varepsilon'^{i-p} (rr'\Psi)^k \cos [j(\theta - \theta')]}{k!p!(i-p)!} \frac{\partial^{i-p}}{\partial a'^{i-p}} \\ &\quad \left[\frac{a'^{-(2k+p+s)}}{2} \frac{d^p}{d\alpha^p} b_{\frac{2k+s}{2}}^{(j)}(\alpha) \right] \end{aligned}$$

3. Derivatives of $B(a, a')$ to order 6

In the following, we have the derivatives of B , with respect to a and a' , calculated by using the derivatives of b as a function of α :

$$\begin{aligned} B(a, a') &= \frac{a'^{-(2k+s)}}{2} b(\alpha) \\ \frac{\partial^i B}{\partial a^i} &= \frac{a'^{-(2k+s+i)}}{2} \frac{d^i b(\alpha)}{d\alpha^i}, \quad \forall i = 1, 2, \dots \\ \frac{\partial B}{\partial a} &= \frac{a'^{-(2k+s)}}{2} \left[-\frac{(2k+s)b}{a'} - \frac{a}{a'^2} \frac{db(\alpha)}{d\alpha} \right] = -\frac{a'^{-(2k+s+1)}}{2} \left[(2k+s)b + \alpha \frac{db(\alpha)}{d\alpha} \right], \\ \frac{\partial}{\partial a} \left(\frac{\partial B}{\partial a'} \right) &= -\frac{a'^{-(2k+s+2)}}{2} \left[(2k+s+1) \frac{db(\alpha)}{d\alpha} + \alpha \frac{d^2 b(\alpha)}{d\alpha^2} \right], \\ \frac{\partial^2 B}{\partial a'^2} &= \frac{a'^{-(2k+s+2)}}{2} \left[(2k+s)(2k+s+1)b + 2(2k+s+1)\alpha \frac{db(\alpha)}{d\alpha} + \alpha^2 \frac{d^2 b(\alpha)}{d\alpha^2} \right] \end{aligned}$$

$$\begin{aligned} \frac{\partial}{\partial a} \left(\frac{\partial^2 B}{\partial a'^2} \right) &= \frac{a'^{-(2k+s+3)}}{2} \left[(2k+s+1)(2k+s+2) \frac{db(\alpha)}{d\alpha} \right. \\ &\quad \left. + 2(2k+s+2)\alpha \frac{d^2b(\alpha)}{d\alpha^2} + \alpha^2 \frac{d^3b(\alpha)}{d\alpha^3} \right] \\ \frac{\partial^2}{\partial a^2} \left(\frac{\partial B}{\partial a'} \right) &= -\frac{a'^{-(2k+s+3)}}{2} \left[(2k+s+2) \frac{d^2b(\alpha)}{d\alpha^2} + \alpha \frac{d^3b(\alpha)}{d\alpha^3} \right] \\ \frac{\partial^3 B}{\partial a'^3} &= -\frac{a'^{-(2k+s+3)}}{2} \left\{ (2k+s)(2k+s+1)(2k+s+2)b \right. \\ &\quad \left. + 3(2k+s+1)(2k+s+2)\alpha \frac{db(\alpha)}{d\alpha} \right. \\ &\quad \left. + 3(2k+s+2)\alpha^2 \frac{d^2b(\alpha)}{d\alpha^2} + \alpha^3 \frac{d^3b(\alpha)}{d\alpha^3} \right\} \\ \frac{\partial^3}{\partial a'^3} \left(\frac{\partial B}{\partial a'} \right) &= -\frac{a'^{-(2k+s+4)}}{2} \left[(2k+s+3) \frac{d^3b(\alpha)}{d\alpha^3} + \alpha \frac{d^4b(\alpha)}{d\alpha^4} \right] \\ \frac{\partial^2}{\partial a^2} \left(\frac{\partial^2 B}{\partial a'^2} \right) &= \frac{a'^{-(2k+s+4)}}{2} \left[(2k+s+2)(2k+s+3) \frac{d^2b(\alpha)}{d\alpha^2} \right. \\ &\quad \left. + 2(2k+s+3)\alpha \frac{d^3b(\alpha)}{d\alpha^3} + \alpha^2 \frac{d^4b(\alpha)}{d\alpha^4} \right], \\ \frac{\partial}{\partial a} \left(\frac{\partial^3 B}{\partial a'^3} \right) &= -\frac{a'^{-(2k+s+4)}}{2} \left\{ (2k+s+1)(2k+s+2)(2k+s+3) \frac{db(\alpha)}{d\alpha} \right. \\ &\quad \left. + 3(2k+s+2)(2k+s+3)\alpha \frac{d^2b(\alpha)}{d\alpha^2} \right. \\ &\quad \left. + 3(2k+s+3)\alpha^2 \frac{d^3b(\alpha)}{d\alpha^3} + \alpha^3 \frac{d^4b(\alpha)}{d\alpha^4} \right\} \\ \frac{\partial^4 B}{\partial a'^4} &= \frac{a'^{-(2k+s+4)}}{2} \left\{ (2k+s)(2k+s+1)(2k+s+2)(2k+s+3)b \right. \\ &\quad \left. + 4(2k+s+1)(2k+s+2)(2k+s+3)\alpha \frac{db(\alpha)}{d\alpha} \right. \\ &\quad \left. + 6(2k+s+2)(2k+s+3)\alpha^2 \frac{d^2b(\alpha)}{d\alpha^2} \right. \\ &\quad \left. + 4(2k+s+3)\alpha^3 \frac{d^3b(\alpha)}{d\alpha^3} + \alpha^4 \frac{d^4b(\alpha)}{d\alpha^4} \right\} \\ \frac{\partial^5 B}{\partial a'^5} &= -\frac{a'^{-(2k+s+5)}}{2} \left\{ (2k+s)(2k+s+1)(2k+s+2)(2k+s+3) \right. \\ &\quad \left. (2k+s+4)b(\alpha) \right. \\ &\quad \left. + 5(2k+s+1)(2k+s+2)(2k+s+3)(2k+s+4)\alpha \frac{db(\alpha)}{d\alpha} \right\} \end{aligned}$$

$$\begin{aligned}
& +10(2k+s+2)(2k+s+3)(2k+s+4)\alpha^2 \frac{d^2b(\alpha)}{d\alpha^2} \\
& +10(2k+s+3)(2k+s+4)\alpha^3 \frac{d^3b(\alpha)}{d\alpha^3} \\
& +5(2k+s+4)\alpha^4 \frac{d^4b(\alpha)}{d\alpha^4} + \alpha^5 \frac{d^5b(\alpha)}{d\alpha^5} \Big\} \\
\frac{\partial}{\partial a} \left(\frac{\partial^4 B}{\partial a'^4} \right) &= \frac{a'^{-(2k+s+5)}}{2} \Big\{ (2k+s+1)(2k+s+2)(2k+s+3) \\
(2k+s+4) \frac{db(\alpha)}{d\alpha} &+ 4(2k+s+2)(2k+s+3)(2k+s+4)\alpha \frac{d^2b(\alpha)}{d\alpha^2} \\
+6(2k+s+3)(2k+s+4)\alpha^2 &\frac{d^3b(\alpha)}{d\alpha^3} \\
+4(2k+s+4)\alpha^3 \frac{d^4b(\alpha)}{d\alpha^4} &+ \alpha^4 \frac{d^5b(\alpha)}{d\alpha^5} \Big\} \\
\frac{\partial^2}{\partial a^2} \left(\frac{\partial^3 B}{\partial a'^3} \right) &= -\frac{a'^{-(2k+s+5)}}{2} \Big\{ (2k+s+2)(2k+s+3)(2k+s+4) \frac{d^2b(\alpha)}{d\alpha^2} \\
+3(2k+s+3)(2k+s+4)\alpha &\frac{d^3b(\alpha)}{d\alpha^3} \\
+3(2k+s+4)\alpha^2 \frac{d^4b(\alpha)}{d\alpha^4} &+ \alpha^3 \frac{d^5b(\alpha)}{d\alpha^5} \Big\} \\
\frac{\partial^3}{\partial a^3} \left(\frac{\partial^2 B}{\partial a'^2} \right) &= \frac{a'^{-(2k+s+5)}}{2} \Big\{ (2k+s+3)(2k+s+4) \frac{d^3b(\alpha)}{d\alpha^3} \\
+2(2k+s+4)\alpha \frac{d^4b(\alpha)}{d\alpha^4} &+ \alpha^2 \frac{d^5b(\alpha)}{d\alpha^5} \Big\} \\
\frac{\partial^4}{\partial a^4} \left(\frac{\partial B}{\partial a'} \right) &= -\frac{a'^{-(2k+s+5)}}{2} \left[(2k+s+4) \frac{d^4b(\alpha)}{d\alpha^4} + \alpha \frac{d^5b(\alpha)}{d\alpha^5} \right] \\
\frac{\partial^6 B}{\partial a'^6} &= \frac{a'^{-(2k+s+6)}}{2} \Big\{ (2k+s)(2k+s+1)(2k+s+2) \\
(2k+s+3)(2k+s+4)(2k+s+5)b(\alpha) & \\
+6(2k+s+1)(2k+s+2)(2k+s+3)(2k+s+4) & \\
(2k+s+5)\alpha \frac{db(\alpha)}{d\alpha} &+ \\
+15(2k+s+2)(2k+s+3)(2k+s+4)(2k+s+5)\alpha^2 &\frac{d^2b(\alpha)}{d\alpha^2} \\
+20(2k+s+3)(2k+s+4)(2k+s+5)\alpha^3 &\frac{d^3b(\alpha)}{d\alpha^3}
\end{aligned}$$

$$\begin{aligned}
 & +15(2k + s + 4)(2k + s + 5)\alpha^4 \frac{d^4 b(\alpha)}{d\alpha^4} + \\
 & +6(2k + s + 5)\alpha^5 \frac{d^5 b(\alpha)}{d\alpha^5} + \alpha^6 \frac{d^6 b(\alpha)}{d\alpha^6} \} \\
 \frac{\partial}{\partial a} \left(\frac{\partial^5 B}{\partial a'^5} \right) &= -\frac{a'^{-(2k+s+6)}}{2} \left\{ (2k + s + 1)(2k + s + 2)(2k + s + 3) \right. \\
 & \quad (2k + s + 4)(2k + s + 5) \frac{db(\alpha)}{d\alpha} \\
 & +5(2k + s + 2)(2k + s + 3)(2k + s + 4)(2k + s + 5)\alpha \frac{d^2 b(\alpha)}{d\alpha^2} \\
 & +10(2k + s + 3)(2k + s + 4)(2k + s + 5)\alpha^2 \frac{d^3 b(\alpha)}{d\alpha^3} \\
 & \quad 10(2k + s + 4)(2k + s + 5)\alpha^3 \frac{d^4 b(\alpha)}{d\alpha^4} \\
 & \left. +5(2k + s + 5)\alpha^4 \frac{d^5 b(\alpha)}{d\alpha^5} + \alpha^5 \frac{d^6 b(\alpha)}{d\alpha^6} \right\} \\
 & \quad \frac{\partial^2}{\partial a'^2} \left(\frac{\partial^4 B}{\partial a'^4} \right) = \\
 \frac{a'^{-(2k+s+6)}}{2} & \left\{ (2k + s + 2)(2k + s + 3)(2k + s + 4)(2k + s + 5) \frac{d^2 b(\alpha)}{d\alpha^2} \right. \\
 & +4[(2k + s + 3)(2k + s + 4)(2k + s + 5)]\alpha \frac{d^3 b(\alpha)}{d\alpha^3} \\
 & + [6(2k + s + 4)(2k + s + 5)]\alpha^2 \frac{d^4 b(\alpha)}{d\alpha^4} \\
 & \left. +4(2k + s + 5)\alpha^3 \frac{d^5 b(\alpha)}{d\alpha^5} + \alpha^4 \frac{d^6 b(\alpha)}{d\alpha^6} \right\} \\
 \frac{\partial^3}{\partial a'^3} \left(\frac{\partial^3 B}{\partial a'^3} \right) &= -\frac{a'^{-(2k+s+6)}}{2} \left\{ (2k + s + 3)(2k + s + 4)(2k + s + 5) \frac{d^3 b(\alpha)}{d\alpha^3} \right. \\
 & +3(2k + s + 4)(2k + s + 5)\alpha \frac{d^4 b(\alpha)}{d\alpha^4} \\
 & \left. +3(2k + s + 5)\alpha^2 \frac{d^5 b(\alpha)}{d\alpha^5} + \alpha^3 \frac{d^6 b(\alpha)}{d\alpha^6} \right\} \\
 \frac{\partial^4}{\partial a'^4} \left(\frac{\partial^2 B}{\partial a'^2} \right) &= \frac{a'^{-(2k+s+6)}}{2} \left\{ [(2k + s + 4)(2k + s + 5) \frac{d^4 b(\alpha)}{d\alpha^4} \right. \\
 & \left. +2(2k + s + 5)] \alpha \frac{d^5 b(\alpha)}{d\alpha^5} + \alpha^2 \frac{d^6 b(\alpha)}{d\alpha^6} \right\} \\
 \frac{\partial^5}{\partial a'^5} \left(\frac{\partial B}{\partial a'} \right) &= -\frac{a'^{-(2k+s+6)}}{2} \left[(2k + s + 5) \frac{d^5 b(\alpha)}{d\alpha^5} + \alpha \frac{d^6 b(\alpha)}{d\alpha^6} \right]
 \end{aligned}$$

We use the following notations:

$$db(j, k + s/2, p, i - p) = \frac{\partial^p}{\partial a^p} \frac{\partial^{i-p}}{\partial a^{i-p}} \left[\frac{a^{i-(2k+s)}}{2} b_{k+\frac{s}{2}}^{(j)}(\alpha) \right]$$

$$b(j, k + \frac{s}{2}, \alpha) \equiv dbal(s, j, k, 0) \equiv B_{k+\frac{s}{2}}^{(j)}(\alpha) = \left[\frac{a^{i-(2k+s)}}{2} b_{k+\frac{s}{2}}^{(j)}(\alpha) \right], \quad k = 0, 1, \dots, 6$$

$$dbal(s, j, k, p) = \frac{\partial^p}{\partial \alpha^p} b_{k+\frac{s}{2}}^{(j)}(\alpha), \quad p = 1, 2, \dots, 6; \quad k = 0, 1, \dots, 6$$

in particular, we have:

$$dbal(s, j, k, 0) \equiv B_{k+\frac{s}{2}}^{(j)}(\alpha) = ,$$

and

$$b\left(j, \frac{1}{2}, \alpha\right) \equiv dbal(s, j, 0, 0) \equiv B_{\frac{s}{2}}^{(j)}(\alpha).$$

4. Concluding remarks

We postulated the most general formulae, for the literal expansion of Δ^{-s} . In analytical planetary theories, s takes the values 1, 3, 5, ... according to the order in planetary masses. Brouwer and Clemence [5], Brown and Shook [4], Abu El Ata and Chapront [6], Kaula [12], Kamel [7, 8], Murray [9], are responsible for low order expansions. Abu El Ata and Chapront eliminated some of the draw backs. Kaula developed methods of construction in terms of standard orbital elements. Murray extended the expansion of the disturbing function by considering $n > 2$, where n is the number of interacting planets. Kamel used the method of differential operators to carry out the expansion of Δ^{-s} , to degree 3, in eccentricity–inclination, and further in terms of Poincare' variables L, λ, H, K, P . Soliman [11] developed the expansion of Δ^{-s} , following the method of W.M. Smart, neglecting terms of degree > 4 in eccentricity–inclination.

In our present research paper, we applied Taylor's theory, in our derivation to acquire the formula of Δ^{-s} . In Part II the final explicit expansion, will be represented by a Poisson trigonometric series of cosines, with polynomial variables a, e, I , and harmonic variables λ, ϖ, Ω , after the implementation of a computer algebraic *MACSYMA* program.

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